

Supporting Information

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SI Text

Phonon Scattering Rate. Fermi's Golden Rule states that transition rate Γ_i^f from an initial state Φ_i to a final state Φ_f under perturbation δH is $\Gamma_i^f = \langle \Phi_f | \delta H | \Phi_i \rangle$. In the present study, we consider two types of perturbation on a harmonic crystalline system—isotope mass disorder and intrinsic third-order lattice anharmonicity.

Within the single mode excitation approximation, the phonon scattering rate due to mass disorder is static (i.e., temperature independent), and is shown as (1):

$$\begin{aligned} \Gamma_{\text{iso}}(\vec{q}, i) &= \frac{2\pi}{\hbar} \left(\frac{\hbar\omega(\vec{q}, i)}{2} \right)^2 \frac{V_o}{8\pi^3} \\ &\times \sum_{i=1}^{3n} \int_{BZ} d\vec{q}' \left[(\delta(\hbar\omega(\vec{q}, i) - \hbar\omega(\vec{q}', i')) \right. \\ &\times \left. \sum_k g_k |\vec{e}_k(\vec{q}, i) \cdot \vec{e}_k(\vec{q}', i')|^2 \right], \end{aligned} \quad [\text{S1}]$$

where $g_k = \sum_n f_n(k)(1 - m_n(k)/\bar{m}(k))^2$, \hbar is the Plank constant, $f_n(k)$ is the fraction of n -th isotope of atom k that has mass $m_n(k)$, and \bar{m} is the average mass of atom k . $\omega(\vec{q}, i)$ and $\vec{e}(\vec{q}, i)$ are the phonon frequency and polarization vector of the phonon mode (\vec{q}, i) , respectively.

The phonon scattering rate induced by the third-order lattice anharmonicity is expressed as the following:

$$\begin{aligned} \Gamma_{\text{anh}}(\vec{q}, i) &= \frac{V_o}{8\pi^3} \iint_{BZ} d\vec{q}' d\vec{q}'' \\ &\times \sum_{i'} \sum_{i''} \left\{ ((n(\vec{q}'', i'') - n(\vec{q}', i')) S_{\vec{q}, \vec{q}', \vec{q}''}^{\vec{q}, i'} \right. \\ &\left. - \frac{1}{2} ((n(\vec{q}', i') + n(\vec{q}'', i'') + 1) S_{\vec{q}, \vec{q}', \vec{q}''}^{\vec{q}, i'} \right\}, \end{aligned} \quad [\text{S2}]$$

Where,

1. Tamura S (1983) Isotope scattering of dispersive phonons in Ge. *Phys Rev B* 27:858–866.

$$\begin{aligned} S_{\vec{q}, \vec{q}', \vec{q}''}^{\vec{q}, i'} &= \frac{\hbar^2 \delta(\hbar\omega(\vec{q}, i) + \hbar\omega(\vec{q}', i') - \hbar\omega(\vec{q}'', i'')) \delta_{\vec{q} + \vec{q}' - \vec{q}'', \vec{g}}}{24^2 \pi \omega(\vec{q}, i) \omega(\vec{q}', i') \omega(\vec{q}'', i'')} \\ &\times \left| \sum_{lmn} \sum_{\alpha\beta\gamma} \frac{A_{\alpha\beta\gamma}^{lmn}(\vec{q}, \vec{q}', \vec{q}'') e_{\alpha}^l(\vec{q}', i') e_{\beta}^m(\vec{q}', i') e_{\gamma}^n(\vec{q}'', i'')}{\sqrt{m_l m_m m_n}} \right|^2 \end{aligned}$$

and

$$\begin{aligned} S_{\vec{q}, \vec{q}', \vec{q}''}^{\vec{q}, i'} &= \frac{\hbar^2 \delta(\hbar\omega(\vec{q}, i) + \hbar\omega(\vec{q}', i') - \hbar\omega(\vec{q}'', i'')) \delta_{\vec{q} - \vec{q}' - \vec{q}'', \vec{g}}}{24^2 \pi \omega(\vec{q}, i) \omega(\vec{q}', i') \omega(\vec{q}'', i'')} \\ &\times \left| \sum_{lmn} \sum_{\alpha\beta\gamma} \frac{A_{\alpha\beta\gamma}^{lmn}(\vec{q}, \vec{q}', \vec{q}'') e_{\alpha}^l(\vec{q}, i) e_{\beta}^m(\vec{q}', i') e_{\gamma}^n(\vec{q}'', i'')}{\sqrt{m_l m_m m_n}} \right|^2 \end{aligned}$$

$n(\vec{q}, i) = \frac{1}{e^{\hbar\omega(\vec{q}, i)/k_B T} - 1}$ is the Bose–Einstein distribution function, m is the mass, ω is the phonon frequency, and $A_{\alpha\beta\gamma}^{lmn}(\vec{q}, \vec{q}', \vec{q}'')$ is the third-order lattice anharmonicity tensor in the reciprocal space. We have developed an efficient real-space supercell numerical algorithm to calculate the third-order lattice anharmonicity tensors by using first-principles methods. Computational details and the results of density/pressure dependence of lattice anharmonicity in MgO can be found in our previous study (2).

The overall transition rate of a phonon mode $\Gamma(\vec{q}, i)$ is approximated as the sum of $\Gamma_{\text{iso}}(\vec{q}, i)$ and $\Gamma_{\text{anh}}(\vec{q}, i)$ (Matthiessen's rule) that are evaluated based on first-principles calculated phonon frequencies, phonon polarization vectors as well as the third-order lattice anharmonicity tensors. A $16 \times 16 \times 16$ grid sampling provides numerically converged results for the two-atom unit cell of MgO crystals. Direct calculation of phonon lifetime is carried out only for those phonon modes at the irreducible \vec{q} -point in the first Brillouin zone, and the rest are reconstructed with group theory based on the crystal symmetry.

2. Tang X, Dong J (2009) Pressure dependence of harmonic and anharmonic lattice dynamics in MgO: A first-principles calculation and implications for lattice thermal conductivity. *Phys Earth Planet In* 174:33–38.