Fold-change detection and scale invariance of cell–cell signaling in social amoeba

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Cell–cell signaling is subject to variability in the extracellular volume, cell number, and dilution that potentially increase uncertainty in the absolute concentrations of the extracellular signaling molecules. To direct cell aggregation, the social amoeba Dicyostelium discoideum collectively give rise to oscillations and waves of cyclic adenosine 3′,5′-monophosphate (cAMP) under a wide range of cell density. To date, the systems-level mechanism underlying the robustness is unclear. By using quantitative live-cell imaging, here we show that the magnitude of the cAMP relay response of individual cells is determined by fold change in the extracellular cAMP concentrations. The range of cell density and exogenous cAMP concentrations that support oscillations at the population level agrees well with conditions that support a large fold-change–dependent response at the single-cell level. Mathematical analysis suggests that invariance of the oscillations to density transformation is a natural outcome of combining secrete-and-sense systems with a fold-change detection mechanism.

Cell–cell signaling lies at the basis of development and maintenance of multicellular forms of life. Extracellular signals are often subject to greater fluctuations in the size of extracellular space and the number of cells (Fig. L4), not to mention nonspecific binding to other molecules, degradation, and dilution. These factors introduce an uncertainty to the detectable number of extracellular ligand molecules, thus posing a threat to the fidelity of cell–cell communication. One of the means by which cells could cope with such uncertainties is to base their behavioral decisions on temporal changes in the extracellular signals. Persistent stimuli are often ignored while their changes in time elicit transient responses—a property collectively called adaptation (1–3). Recent studies have highlighted cellular response whose magnitude appears to be dictated by the fold change in the input stimuli—a property referred to as “fold-change detection” (FCD) (4, 5). In bacterial chemotaxis, cells respond adaptively to a fold change in chemotacticant concentration (6) so that their search patterns depend only on the spatial profiles of the chemotacticant irrespective of its absolute level. Fold-change dependence is also implied in eukaryotic chemotactic response (7, 8) as well as cell fate control and gene regulation in Xenopus embryo (9), Drosophila imaginal disk (10), and mammalian cells (11). These studies have shed light on the role of FCD for a simple unidirectional signal transduction from an extracellular ligand-receptor interaction (input) to a cellular response (output). However, cell-cell signaling and multicellular systems as a whole often use secretion and sensing of the same molecules (12), whereby the output is fed back to the responding cell itself in addition to the neighboring cells, thus forming a complex bidirectional signal transduction system. The consequence of equipping such systems with an adaptive response and FCD is so far unaddressed.

Oscillations of extracellular cAMP that dictate aggregation of the social amoebae Dicyostelium discoideum is a prime example of robust collective behaviors in cell populations. Under starvation, cells synthesize and secrete cAMP, which stimulates other cells in the vicinity to induce further synthesis and secretion of cAMP—a process called “cAMP relay” (13). After prolonged exposure to cAMP, the rise in extracellular cAMP level ceases due to inactivation of adenyl cyclase (14). As extracellular cAMP level is lowered by degradation, the cells exit from the state of reduced responsivity over the course of several minutes (15, 16), and hence the extracellular cAMP level once again starts to elevate. This tendency for the extracellular cAMP level to rise when it is lowered, and to be lowered when it is raised, essentially renders extracellular cAMP level unstable and oscillatory. The emerging oscillatory waves of extracellular cAMP in the cell population provide a temporal guidance cue for directional cell migration (17, 18). Both cAMP oscillations and cell aggregation are known to occur at a wide range of cell densities spanning at least two to three orders of magnitude (19–22). Cells lacking the cAMP-synthesizing enzyme adenyl cyclase ACA when ectopically forced to differentiate cannot form aggregates unless they are allowed to randomly collide and form cell clusters at high cell densities (23), indicating that robust cell aggregation depends on intracellular cAMP signaling. Recent live-cell imaging studies have elucidated an input-output relation of the cAMP relay response at the single-cell level resolution (16, 24–26). Although these analyses have implied a role played by the stochastic cAMP relay response at low basal concentrations of extracellular cAMP (~100 pM) in initiating the synchronized pulses, how they could take place robustly in a wide range of cell densities is yet poorly understood.

Results
Density Dependence of the Oscillations. To gain insights on the concentration range of intracellular and extracellular cAMP that

Significance
Recent works have hinted at an ability of cells to respond in the exact same manner to a fold change in the input stimulus. The property is thought to allow cells to function properly regardless of changes in the absolute concentrations of signaling molecules. Despite its general importance, however, evidence has remained scarce. The present work demonstrated that, in the social amoeba Dicyostelium, a response to cell–cell communication molecules is fold-change dependent and that this property is tightly linked to the condition that allows them to oscillate collectively, and thus to organize into a multicellular form. Such properties may be of importance for robustness of other developmental systems where oscillatory signaling plays a pivotal role in defining multicellular organization.

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supports oscillations, we revisited the cAMP dynamics quantitatively by live-cell imaging of cells expressing the cAMP probe Epac1camps (24) combined with a corrected quantification of the fluorescence resonance energy transfer (FRET) efficiency (27) (“FRET index”; Fig. S1 and SI Materials and Methods). Collective cAMP oscillations occur spontaneously in cell populations continuously perfused with buffer solution at a moderate flow rate (1.5 mL/min), where extracellular cAMP breakdown by secreted phosphodiesterase (PDE) in vivo is effectively emulated by dilution (24, 28). First, we tested the effect of adding cAMP in the perfusion flow (hereafter refer to as “exogenous cAMP”), which adds on top of endogenously synthesized and secreted cAMP present in the extracellular space. For 0.3 nM exogenous cAMP, the oscillations persisted at a relatively high cell density (1/2 ML) (Fig. 1A, Left Upper), whereas almost no clear effect was observed at a high cell density (1/2 ML). (C) A phase diagram of the mean pulse frequency as a function of cell density and exogenous cAMP. At high cell densities, the oscillations were relatively insensitive to exogenous cAMP. Squares represent three to six independent measurements of a 2−h duration (n = 4, 4, 3, and 3 for 1/2 ML from Left to Right; n = 8, 4, 3, and 3 for 1/8 ML; n = 4, 3, and 3 for 1/128 ML). (D) Representative time series of the FRET index (Epac1camps/AAX4 cells) obtained from individual cells oscillating in a population of 1/8 monolayer (ML) (Left) and 1/2 ML cell density; perfusion rate, 1.5 mL/min. Colors indicate different cells. (E) The average of minimum (Left) and maximum (Right) FRET index. Error bars represent SD (n = 21 for 1/8 ML and n = 19 for 1/8 ML).

**Rescaled Response Sensitivity.** What exact concentration and magnitude change in the extracellular cAMP are required to sustain the collective oscillations? Before and during the early stage of cell aggregation, the basal level of extracellular cAMP slowly builds up and is estimated to reach around 10 nM or less (29, 30), which is near the Kd of the membrane-bound G-protein-coupled receptor CAR1. To see the effect of background concentrations of extracellular cAMP in this range, we measured single-cell level cAMP relay responses in well-isolated cells (<10−7 ML) (24) that were first primed with a fixed concentration of cAMP (0.1, 1, 3, and 10 nM) for about 40 min to let the initial response attenuate before elevating the cAMP level in a step-like manner (Materials and Methods). Note that, in contrast to the population-level oscillation experiments described in the previous section, here, care was taken so that cAMP secreted by the cells becomes negligible (24); extracellular buffer was rapidly diluted by fast perfusion (4–8 mL/s; Materials and Methods), cells were plated at low cell densities (see Materials and Methods), and a cell was chosen for observation after confirming that there is no other cell within the surrounding area of ∼0.3-mm radius. Fig. 24 shows reference time courses of the cAMP relay response to a step increase in the extracellular cAMP concentration from 0 to 3 nM, a nonprimed stimulus condition that elicits a maximum amplitude response as measured by the FRET signal (24). The level of cytosolic cAMP reached the maximum at 2−3 min after the stimulus application followed by secondary peaks that slowly attenuated in close agreement with a previous study (24). Under primed conditions, the response to 3 nM on top of 1 nM extracellular cAMP was markedly reduced in both initial and secondary peak amplitude (Fig. 2B, Left), whereas elevation from 1 to 10 nM induced a large magnitude response (Fig. 2B, Right) comparable to that observed under nonpriming conditions. Next, we quantified the initial peak amplitude 2−3 min after the stimulus increase. This serves as the primary measure of the strength of the relay response underlying the collective oscillations, because the secondary peaks were not observed when the level of extracellular cAMP was elevated only for a few minutes.
The cAMP relay response at the single-cell level obeys fold-change detection. (A and B) The concentration of perfused cAMP was changed from 0 to 3 nM (A) to 10 nM (B). The dotted lines indicate the time of stimulus switch. Colors indicate different cells. (C and D) The peak amplitude (C) and the response time (D) of the response; priming concentrations, 0.1 nM (blue), 1 nM (green), 3 nM (red), and 10 nM (magenta) cAMP. The vertical axis: the maximum FRET index during the 4 min after the input change (n = 7). The dotted lines indicate the time of stimulus switch. Colors indicate different cells. (E and F) The response amplitude (E) and the response time (F) of the response; priming concentrations, 0.1 nM (blue), 1 nM (green), 3 nM (red), and 10 nM (magenta) cAMP. The vertical axis: the maximum FRET index during the 4 min after the input change (n = 7). Dots and error bars indicate the mean and SD (n = 5). The results indicate that the fold-change dependence is a single-cell level property.

(24) (Fig. S2A) as is the case during the collective oscillations. Fig. 2C plots the index as a function of fold increase in extracellular cAMP concentration for four priming concentrations tested. For priming concentrations below 10 nM (0.1, 1, and 3 nM), the plots collapsed on a single curve, indicating that, within this concentration range, the single major determinant of the peak amplitude of the cAMP relay response is the fold change in the extracellular cAMP concentration. The peak FRET index observed during the relay response was half-maximal at about a 3-fold increase in the extracellular cAMP concentration and levels off at about a 10-fold increase. The time at which the initial peak reached the maximum value showed almost no dependency on the background level of extracellular cAMP (Fig. 2D). Moreover, the amplitude of the response to a fold change of 10 was more or less constant between 0.1 and 10 nM background concentrations (Fig. 2E), suggesting that the rescaling property holds for two orders of magnitude. The same rescaling property was also observed when the step-stimulus was repeated (Fig. S2A).

One should note that an apparent fold-change dependence in the population-averaged response (Fig. 2C) should not be readily equated with a fold-change detection at the single-cell level. For example, cell–cell variability in the response sensitivity alone in the absence of FCD could deceptively bring about a similar response curve after averaging (Fig. 2F, Left). In principle, whether or not rescaling holds at the level of individual cell (Fig. 2F, Right) should be tested by measuring the response to fold-change stimulus numerous times in single cells. However, this is not feasible due to phototoxicity and also because cells differentiate in the timescale of hours. Instead, we measured the response to two incremental steps with a fold change of 5 (0.3, 1.5, and 7.5 nM cAMP). Fig. 2G shows representative time series of single-cell level responses. Those cells that show weak response in the initial increment also tend to show weak response in the second increment. Likewise, cells with large response were consistent in two incremental stimuli. Despite a large cell–cell variability in the amplitude, the variability was spread along a linear slope in the scatter plot (Fig. 2H). In other words, the amplitudes of the response to the consecutive inputs were highly correlated in each individual cells. The results indicate that the fold-change dependence is a single-cell level property.

During cell aggregation of Dicyostelium discoideum, synchronous emission of cAMP gives rise to periodic traveling waves of extracellular cAMP, and thus the level of extracellular cAMP is oscillating (29). In this respect, the spatially uniform and temporally abrupt stimulus used in the above section deviates from the naturally occurring one. In general, a response to a step stimulus may not necessarily have the same propensity as for a slowly varying stimulus. Moreover, directional migration is induced by a traveling-wave stimulus (17, 18) and thus may affect the cAMP relay owing to a large overlap in the signal transduction network (16, 31, 32). To test relevance of the fold-change response property for natural cAMP wave, we used a microfluidics lighthouse (33) —a gradient-generating platform capable of delivering traveling-wave stimulus of various amplitude, frequencies, and speed. Using this device, we applied 6-min period cAMP stimuli to cells in the region of interest (Fig. 3A; dotted square) in the form of spatially symmetric bell-shaped gradient that measures ~500 μm from the trough to the peak and traversed at a speed of ~2–3 × 10–4 μm/min (Fig. 3A and B), which closely follows the estimated parameters of natural cAMP waves (34). The flow is slow and does not bias directionality of cell migration (33), and there is no apparent shear-induced effect on the level of cytosolic cAMP (Fig. 3C). Fig. 3 D–F shows representative responses of well-isolated cells to repetitive wave stimulus with fold increase/decrease of 30 for background concentration of 0.1, 1.0, and 3.0 nM extracellular cAMP. In all three conditions, albeit cell–cell variability, the concentration of cytosolic cAMP in individual cells oscillated in phase with that of extracellular cAMP whose peak time showed delay of ~10–40 s. The peak FRET index was on average almost constant under these conditions (Fig. 3G), confirming that the magnitude of single-cell level response to traveling-wave stimulus is fold-change dependent.

The cAMP relay response is mediated by complex signaling cascades downstream of G-protein–coupled receptors that are yet to be
fully revealed (31, 35). Binding of extracellular cAMP to the receptor activates small GTPase Ras, which then activates phosphoinositide 3-kinase (PI3K) to elevate phosphatidylinositol-(3,4,5)-triphosphate (PIP3) at the plasma membrane (36). PIP3 serves as a docking site for a pleckstrin homology (PH) domain-containing protein cytosolic regulator of adenylyl cyclase (CRAC) that is essential for activation of adenylyl cyclase ACA (37, 38). These activation processes are counteracted by deactivation of Ras by RasGAP (37, 38). The orientation of the stimulus flow was varied at the rate of 0.34°/s anticlockwise from θ = 0° to 120°, thereby allowing a bell-shaped gradient to traverse the observation area (dotted square). For the remaining region from θ = 120° to 360°, the rotation was fast-forwarded by switching the direction discretely in two steps from 120° to 240°, and then 240° to 360° (33). Alexa Fluor 594 was used in the stimulus solution to estimate the spread of cAMP. For each run of experimental observation, the rotation was cycled four to six times. (B) Merged confocal images of the stimulus profile (red; Alexa) and Epac1camps/AX4cells (yellow; YFP (I<sub>saximp</sub> channel) during a wave passage. (Scale bar, 50 μm.) (C-F) cAMP response to the wave stimulus. cAMP concentrations: mock (C; n = 5 cells), 0.1–3 nM (D; n = 12 cells), 1–30 nM (E; n = 10 cells), and 3–90 nM (F; n = 6 cells). Temporal profiles of the wave stimulus (Upper) and individual cell responses (Lower; graded blue green colors indicate different cells). Buffer solution containing Alexa Fluor 594 without cAMP was used for the mock control (C). cAMP concentrations were estimated from the fluorescence intensity of the Alexa dye (D–F). (G) The average peak FRET index. Error bars indicate SDs.

Scale Invariance of FCD-Based Signaling. How are the input–output properties revealed at the single-cell level related to the oscillatory behavior at the population level? Let us formulate a secrete-and-sense system where the output “y” is determined by the input signal “z”; conversely, the input z is determined by the response y. To this end, we shall first describe a ligand-induced response by a two-variable system \( \dot{x} = f(x,y,z) \), \( \dot{y} = g(x,y,z) \), where intracellular and extracellular cAMP are represented by y and z, respectively (Fig. 5A). Here, x is an internal variable required for adaptation. Mathematically, the system obeys FCD when the functions f and g satisfy a condition:

\[
\frac{\partial g(p)}{\partial p} = g(x,y,z)
\]

with an appropriate transformation function \( \phi \) and an arbitrary scaling factor \( p > 0 \) (SI Equations). Under this constraint, the variable y responds in an identical manner to two incremental p-fold increases in the input z (Fig. 5B). By introducing a third equation that

\[ 0.1 \text{ and } 1 \text{ nM priming concentrations of cAMP, the degree of membrane translocation of PHCRAC-RFP exhibited a fold-change dependency (Fig. 4D). The time it took for the response to reach its peak also appeared almost identical (Fig. 4E) with a slight decline at a higher fold stimuli as has been noted for Ras activation (7). Rescaling of the response sensitivity was observed in a somewhat narrower range of the background cAMP concentrations (0.1–1 nM) (Fig. 4F) compared with that of the cAMP relay response (0.1–10 nM) (Fig. 2E). On average, there was decrease in the degree of PHCRAC-RFP translocation at priming concentrations higher than 1 nM. These results suggest that information regarding fold change in the extracellular cAMP concentration is encoded and transduced, at least partially, by the degree of membrane translocation of CRAC.

Materials and Methods

Binding of extracellular cAMP to the receptor activates small GTPase Ras, which then activates phosphoinositide 3-kinase (PI3K) to elevate phosphatidylinositol-(3,4,5)-triphosphate (PIP3) at the plasma membrane (36). PIP3 serves as a docking site for a pleckstrin homology (PH) domain-containing protein cytosolic regulator of adenylyl cyclase (CRAC) that is essential for the activation of adenylyl cyclase ACA (37, 38). These activation processes are counteracted by deactivation of Ras by RasGAP (37, 38).
describes the dynamics of extracellular cAMP $z$ in a perfusion chamber (13, 24), we arrive at a closed-loop system (Fig. 5C):

$$
\begin{align*}
\dot{x} &= f(x, y, z) \\
\dot{y} &= g(x, y, z) \\
\dot{z} &= \rho k y - \gamma z,
\end{align*}
$$

where $\rho$, $k$, and $\gamma$ represent the cell density, the secretion rate, and the dilution rate, respectively. The extracellular cAMP $z$ affects the dynamics of the cells—that is, the subsystem $x$ and $y$; hence the system describes “sensing” and secretion of the same molecules (12, 44).

From the above mathematical formulation, it is readily shown that Eq. 2 are invariant under a transformation; that is, $(x, y, z, \rho) \rightarrow (\phi(p, x), y, p z, p p)$ (see Supporting Information for proof). The scale invariance ensures that, although extracellular cAMP $z$ scales with the cell density, the intracellular cAMP $y$, regardless of whether it is a fixed constant or dynamically changing variable, remains unaffected by density transmutation $\rho \rightarrow p p$. Because the invariance derives simply by requiring the FCD condition (Eq. 1) for the response dynamics (Eq. 2; equations for $x$ and $y$), and combining them with the secretion/dilution dynamics (Eq. 2; the equation for $z$), it can be said that the robustness of the cAMP oscillations to variation in cell density is a natural outcome of the observed response-rescaling property of the constituent cells. Note that, for the sake of clarity, the above derivation assumed that the response functions $f(x, y, z)$ and $g(x, y, z)$ are the same in all cells. The conclusion holds when there is cell–cell variability (SI Equations; Eqs. S6 and S7).

**Scale Invariance of the cAMP Oscillations.** To obtain networks that exhibit adaptive response that obeys FCD, which could then be combined with the equation for $z$ (Eq. 2 and Fig. 5C), a computational search for minimal networks with two nodes and three links (3, 45, 46) was performed (Fig. S4A–C). Four networks were identified: two “incoherent feedforward loop” type (Fig. S4D) and two negative-feedback loops type (Fig. S4E), which are consistent with recent analytical search of a topology space (46). The basic network models could be slightly modified to take into account the nonlinear input–output relation of the cAMP relay response (Fig. 2C) while preserving the network topology and the FCD property (SI Equations). Although details of the response can be compared to further constrain the network topology (26, 47–49), the present aim is rather to understand the general outcome of incorporating FCD in a secrete-and-sense system. Qualitative features detailed below were conserved in all four selected network types (Fig. S5). For brevity of presentation, a feedforward-type network,

$$
\begin{align*}
\dot{x} &= x + K x \\
\dot{y} &= \frac{z^n}{\lambda} - y \\
\dot{z} &= \rho k y - \gamma z,
\end{align*}
$$

is described below (SI Equations, model [A’]). Parameters $n$, $\lambda$, and $K$ are positive constants. The feedforward network departs receptor-mediated signals that induce both fast activation and slow terminating reactions (1) and is compatible with adaptation and deadaptation of adenyl cyclase that depend strictly on elevation and lowering of the extracellular cAMP level, respectively. Earlier mathematical models of cAMP relay (50, 51) as well as chemotaxis models that describe conversion between GTP-bound and GDP-bound form of Ras (42) also have this network topology.

With an appropriate choice of parameter values, the system (Eq. 3) exhibited oscillations (Fig. S5). The oscillatory solution in the present model (Eq. 3), especially the level of intracellular...
cAMP, y, is insensitive to the value of ρ (Fig. 5B) owing to the scale invariance resulting from the FCD property. This feature is in marked contrast with earlier models (50, 51) where oscillations were observed only in a narrow range of cell densities (< one order of magnitude) irrespective of the dilution rate γ [for model (50), see Fig. S6]. With further inclusion of lower detection limit by replacing z by z + δ on the right-hand side of the equations for x and y (Eq. 3), where δ is a positive constant parameter (SI Equations), we see that the lower limit of cell density that supports collective oscillations becomes a monotonically increasing function of the dilution rate (Fig. 5E)—in agreement with earlier experimental data (figure 2B in ref. 24).

The most important suggestion of the present model is that the range of extracellular cAMP level that confers fully periodic collective oscillations should be dictated by the operation range of FCD. Such a view is supported from the estimates of the extracellular cAMP concentration at various cell densities. Under typical experimental parameters, owing to the high dilution rate (γ ∼ 6 min⁻¹ in the present study) compared with the timescale of intracellular cAMP dynamics (∼1/7 min⁻¹), the concentration of extracellular cAMP (z) can be estimated by applying quasi-steady-state approximation (SI Equations). At cell densities of 1/2 and 1/8 ML, where population of cells exhibited cAMP oscillations at a constant high frequency (Fig. 1C) (24), the estimated basal concentrations of extracellular cAMP are 0.72 nM for 1/2 ML and 0.18 nM for 1/8 ML, which are well within the FCD range (Fig. 2C; ≥0.1 nM). At 1/32 and 1/128 ML where populations showed sporadic or no oscillation, the estimated extracellular cAMP are 0.05 nM for 1/32 ML and 0.01 nM for 1/128 ML, which lie at the border or outside the FCD range (Fig. 2E; <0.1 nM). Thus, the conditions that support full-amplitude ∼7-min period oscillations agree well with the FCD range.

The critical relation between FCD and the oscillatory conditions is also vindicated by the suppression of oscillations in the presence of exogenous cAMP and/or at low cell densities as demonstrated experimentally (Fig. 1 B and C) as well as in the numerical simulations (Fig. 5F). Here, the model equations were extended to include the influx of exogenous cAMP (z₀) by rewriting the equation for z as dz/dt = ρky − γ(z − z₀), which now violates the condition for scale invariance. The phase diagram in the parameter space (z₀, ρ) obtained numerically shows a linear slope that demarcates the oscillatory and nonoscillatory domains (Fig. 5G), which closely resembles the phase diagram obtained experimentally (Fig. 1C). The feature of the phase diagram can be understood by the fold-change responsiveness of the constituent cells as follows. Under no exogenous cAMP, the estimated fold change in the extracellular cAMP level during the oscillations is zpeak(ρ)/zbasal(ρ), where zpeak and zbasal are the peak and basal concentrations of extracellular cAMP, respectively, both of which increase as a function of cell density (SI Equations). Conversely, in the presence of exogenous cAMP, the estimate would be AS(ρ, z₀) = zpeak(ρ) / z₀(1/ρ + z₀). Because AS is a decreasing function of z₀ and increasing function of ρ, at higher concentrations of exogenous cAMP and/or at lower cell densities, the response becomes small and less likely to sustain oscillations. Our single-cell data (Fig. 2C) indicate that the change in the level of extracellular cAMP should be greater than ∼10-fold.

Fig. 5. Density robustness is a basic property of a secrete-and-sense system consisting of FCD elements. (A and B) A reaction scheme (A, “isolated Cell”) and the characteristic behaviors (B) of an FCD system. The time evolution of y and x are described by functions f and g that together constitutes a FCD system (A). Simulated time course of x and y (B, Lower) to incremental increase in z (B, Upper). (C and D) A Schematic diagram of a closed-loop secrete-and-sense system (C, “Communicating cells”) and results from model simulations (Eq. 3) for low (Upper) and high (Lower) cell densities. Here, x and y follow the same reaction scheme as in A except that y is secreted, and thus transformed to z to further stimulate the cells. (E) A phase diagram of the system (Eq. 3) with a lower detection limit. Oscillatory (red squares) and static (blue plus sign (+)) states, respectively. The parameters were r = 1.5, n = 2, K = 4, κ = 2, and γ = 3, unless otherwise indicated.
otherwise, the response amplitude was markedly diminished, and there were no population-level oscillations. Consistent with such a view, the values of $A_r$ for conditions that supported the oscillations were always higher than those in the nonoscillatory conditions and the border between the two domains was at $A_r \approx 10$ (Table S1). These results suggest that the secreting-and-sense system that consists of FCD elements captures the essence of the density dependence of the oscillations.

**Discussion**

The present study demonstrated that the cAMP relay response is adaptive and fold-change dependent at low extracellular cAMP concentrations ($\leq 10$ nM). Based on the observed input–output relation, we proposed a model that consists of elements that obey FCD that well describes the cell density and extracellular cAMP conditions that support the population-level oscillations. Recently, an alternative model based on the Fitzhugh–Nagumo equations was proposed (26). There, it is thought that population-level oscillations emerge as a result of a stochastic threshold-like response to fluctuating levels of extracellular cAMP. Such a view contrasts markedly with the model proposed here where a positive-feedback loop between intracellular and extracellular cAMP and adaptation of the response give rise to oscillations—a deterministic scheme that shares its core idea with other earlier models (50). The appearance of frequency selectivity and a refractory period (26) do not suffice as a criteria for model selection as these behaviors were also reproduced in the current adaptation-based model (Fig. S2 B and C and SI Discussion). Experimentally, the obtained dose–response (Fig. 2C) had no clear threshold but rather increased gradually—a feature not in support of an excitatory noise-driven mechanism. Moreover, according to the Fitzhugh–Nagumo-type mechanism, inclusion of exogenous cAMP should give rise to more pulses of cAMP by increasing the chance to cross the response threshold. Our observations were quite the opposite. The oscillations were suppressed by exogenous cAMP (Fig. 1 B and C). Despite these caveats, we should note that the Fitzhugh–Nagumo-based model captures the sustained (i.e., nonadaptive) oscillatory response behavior (26) that occurs under prolonged stimulus in the non-FCD range ($> 10$ nM cAMP) (24). These oscillations are driven intracellularly without the need for clearing of extracellular cAMP—a feature in favor of the Fitzhugh–Nagumo-type scheme. For the population-level oscillations, in contrast, a strict dependency on extracellular cAMP clearing has been demonstrated by a null-mutant of extracellular phosphodiesterase that cannot oscillate unless it is cleared of extracellular cAMP by perfusion (28). It may be that different cAMP receptor forms (52) or receptor types (53) with different binding constants transduce FCD and non-FCD signals.

The present analysis and the proposed theoretical framework suggest that, when an adaptive response governed by FCD is incorporated into a “secrete-and-sense” circuitry (12, 44), the resulting cell–cell communication is robust to cell density changes due to the scale-invariant property of the system equations. In the present analysis, the extracellular signaling molecule $z$ is described by a relatively simple equation owing to the first-order secretion kinetics (13) and the constant dilution rate realized by perfusion. In addition, density robustness can also be realized in more complex situations, for example, nonlinear secretion kinetics, a density-dependent degradation rate, and spatial homogeneity (SI Equations, Fig. S7, and SI Discussion). Although the FCD mechanism could explain the ability of Dictyostelium cells to execute aggregating behavior at a wide range of cell density (21, 22), the present findings do not rule out other mechanisms that may be at work in conjunction. cAMP-dependent regulation of a gene encoding extracellular PDE can provide robustness (28), in this case by keeping the level of extracellular cAMP within a desirable range. Cell–cell heterogeneity in the response sensitivity could also render oscillations robust to cell density change (Fig. 2F, Left). We should note, however, that FCD at the single cell–level (Fig. 2H) alleviates the requirement for averaging over many cells, which could be critical in nature where averaging is limited by diffusion. Given the ubiquity of adaptation and secrete-and-sense circuits (12), FCD-based robustness may have a wide connotation in other multicellular phenomena. A case in point is the collective oscillations of NADH in yeast cells suspension whose amplitude and frequency are conserved for three orders of magnitude in cell density (54). In higher organisms, morphogen fields are dynamic and their temporal changes appear to be read out to regulate cell fates (10, 55). Traveling waves are also prevalent in embryonic development (56, 57). It would be interesting to explore whether the same mathematical principle applies in these and other systems.

**Materials and Methods**

**Strains and Cell Culture.** *Dictyostelium discoideum* AX4 cells expressing Epac1camps (Epac1camps/AX4) (24), PH domain of CRAC monomer fused to monomeric red fluorescent protein (mRFP) (PHCRAC-RFP/AX4) (43), and both PHCRAC-RFP and Epac1camps (Epac1camps–PHCRAC-RFP/AX4) were used. For coexpression, Epac1camps/AX4 cells were transformed with PHCRAC-RFP expression vector by electroporation. A clonal isolate showing normal development with relatively bright fluorescence for both probes was used for the analysis. Cells were grown at 22 °C in PS medium (58) with 10 μg/mL G418 and/or 60 μg/mL hygromycin B, where appropriate. Typically, 30-MLC cell culture was shaken in an Erlenmeyer flask (250 mL, Belco) at 155 rpm (Taitec; BRA3-FL). Cells were propagated between 10^6 cells per mL. For live-cell imaging, cells were removed from growth medium by centrifugation for 3 min at 700 × g followed by resuspension in fresh development buffer (DB) (10 mM K/Na phosphate buffer, 1 mM CaCl2, 2 mM MgCl2; pH 6.5). The washing step was repeated twice. The cells were resuspended in –1 mL DB at –2 × 10^6 cells per mL and allowed to differentiate for 4.5 h in a centrifuge tube shaken at 155 rpm (Taitec; BRA3-FL).

**Perfusion and Live-Cell Microscopy.** Starved cells were plated on a glass-bottom dish at a cell density below 1 × 10^6 cells per cm^2, and a perfusion chamber was constructed using an insert (RC-37F; Warner Instruments) as previously described (24), with the exception of results shown in Fig. 3, which used the lighthouse device (33) (SI Materials and Methods). After cells have attached to the bottom of the chamber, DB was perfused to remove effects from past stimuli at the rate of 4–8 mL/min, unless otherwise noted. For the cAMP relay study, the flow rate was kept at 8 mL/min for longer than 1 min before and after stepwise changes in the cAMP concentrations. At the given flow rate, one-half of the solution in the chamber was exchanged within 10 s as tested by monitoring fluorescein solution under a microscope. When there is no need to change the stimulus level, the flow rate was reduced to 4 mL/min. For observations of population-level cAMP oscillations, the flow rate was reduced to 1.5 mL/min.

Epifluorescent live-cell imaging of Epac1camps/AX4 cells was performed essentially as described in the previous work (24) using the same inverted microscopy system. In the present study, images were taken at 10–15 s intervals for total duration of ~0.5–2 h. In addition, an image from the YFP channel was acquired to correct for cross-excitation immediately before time-lapse acquisition (SI Materials and Methods). To this end, cells were exposed for 30 ms with 495-nm light using an excitation filter (BP490-500HQ; Olympus), which was further attenuated by 6% using neutral density filters. The emitted light was separated from the excitation by a dichroic mirror (DM505; Olympus).

Confocal imaging of PHCRAC-RFP/AX4 and PHCRAC-RFP Epac1/AX4 was performed using an inverted microscope (IX-81; Olympus) equipped with an automated stage (BIOS-215ST; Sigma Koki), a confocal scanning unit (CSU-X1; Yokogawa), and optical shutters (LS62M2; Uniblitz Electronics). An oil-immersion objective lens (60× PlanApo N, N.A. 1.42, or 20× UPlanSApo N, N.A. 0.85; Olympus) was used. The stage and the shutters were controlled by a stage controller (FC-101G; Sigma Koki) and a shutter driver (VMM-D3; Uniblitz Electronic), respectively. The device control and data acquisition were conducted using the Metamorph software (Molecular Devices). A 445-nm laser (40 mW; Vortran Laser Technology) and a 561-nm laser (25 mW; Melles Griot) were used as light sources. For imaging PHCRAC-RFP, cells were exposed to 561-nm laser for 30 ms with a neutral density filter. The fluorescent light was passed through an emission filter (BA575I; Olympus). For Epac1camps-expressing cells, 445-nm laser attenuated by 50% using neutral density filters was applied for 50 ms. A set
of multiband dichroic mirror and filters (CFP/YFP/HPRed-3X3M-A; Semrock) was used except for the CFP excitation filter (BP425-445HQ; Olympus). Sixteen-bit 512 × 512 pixel images were captured at 3 -to- 10-s interval using an EMCCD camera (Evolve 512; Photometric). All image acquisition was performed at 22 °C. Data were stored in tagged image file format (TIFF) files.

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