

## SUPPORTING INFORMATION

The efficiency of direct interception by a single fiber was computed as (12)

$$E_R = [2E_* \ln E_* - E_* + E_*^{-1}] / \Lambda, \quad [\text{S1}]$$

where

$$E_* = 1 + \frac{d_p}{d}, \quad [\text{S2}]$$

$$\Lambda = 1 - 2 \ln \tau + \frac{\tau^2}{6} - \frac{\tau^4}{144} + \frac{\tau^6}{1080}, \quad [\text{S3}]$$

$$\tau = \pi d \sqrt{W^2 + L^2} / (WL). \quad [\text{S4}]$$

$W$  and  $L$  are width and length of the mesh opening, respectively, and  $d$  is the fiber diameter.

Direct interception efficiency reaches 100% when particle size is equal to mesh width. Although the model assumes  $d_p \ll d$ , the above expressions are still realistic as validated by several other studies where particles are larger than the mesh diameter (12, 17).

The efficiency of diffusional deposition by a single fiber is (12):

$$E_D = 3.7 \Lambda^{-1/3} Pe^{-2/3} + 0.62 Pe^{-1}, \quad [\text{S5}]$$

where  $Pe = dU/D$  is the Peclet number and  $D$  is the diffusion coefficient of the particles. For

colloidal particles or non-motile organisms, diffusion arises from Brownian motion and  $D =$

$kT/(3\pi\rho\nu d_p)$ , where  $k = 1.38 \times 10^{-23} \text{ m}^2 \text{ kg s}^{-2} \text{ K}^{-1}$  is Boltzmann's constant and  $T$  is temperature.

For motile microorganisms diffusivity results from random motility, increasing with swimming

speed, and  $D$  was computed using the semi-empirical results of Visser and Kiørboe ( $D = 2.8d_p$

$^{1.71}$ ,  $D$  in  $\text{cm}^2 \text{ s}^{-1}$  and  $d_p$  in  $\text{cm}$ , ref. 38).

Finally, the total efficiency of a rectangular filter, which was used in the calculations, is

(12)

$$E = \frac{(E_R + E_D)d}{h_E} \left[ 1 - \frac{(E_R + E_D)d}{W + L} \right], \quad [\text{S6}]$$

where  $h_E = WL/(W+L)$  is the equivalent mesh spacing.